

## Hubbard Model for Hydrogen Atom Lattice

---

Consider a *simple cubic lattice* with a hydrogen atom at each lattice site. Assuming that the only relevant electronic state at the hydrogen site is the  $1s$  state, one can describe the system by a model Hamiltonian

$$\mathcal{H} = -t \sum_{\langle ij \rangle \sigma} (c_{i\sigma}^\dagger c_{j\sigma} + c_{j\sigma}^\dagger c_{i\sigma}) + U \sum_i n_{i\uparrow} n_{i\downarrow} \quad (1)$$

where  $i$  represents the lattice site at  $\mathbf{R}_i$ ,  $\langle ij \rangle$  the pair of nearest-neighbor sites  $\mathbf{R}_j = \mathbf{R}_i + \boldsymbol{\delta}$ , ( $\boldsymbol{\delta}$  being  $\pm a\mathbf{e}_\alpha$  with  $\alpha = x, y, z$ ) and  $\sigma$  corresponds to the spin index. In this model, while the on-site Coulomb interaction  $U$  is independent of the lattice constant, the hopping integral  $t$  can be *assumed* to depend on the lattice constant  $a$ :

$$t(a) = \frac{A}{a^\gamma}$$

with  $\gamma > 0$ . Thus, as  $a$  varies from 0 to  $\infty$ , one can take the limit of either  $U \rightarrow 0$  or  $t \rightarrow 0$ .

### a) Density-of-States:

First, considering the free electron limit, i.e.,  $U \rightarrow 0$ , show that the kinetic part of the Hamiltonian  $\mathcal{H}$  in Eq. (1) can be represented by

$$\mathcal{H}_0 = \sum_{\mathbf{k}\sigma} \varepsilon_{0\mathbf{k}} c_{\mathbf{k}\sigma}^\dagger c_{\mathbf{k}\sigma} \quad (2)$$

where the summation over  $\mathbf{k}$  is limited to the first Brillouin zone. Also, obtain that the density-of-state (DOS) at energy  $\varepsilon$  can be written by

$$D_0(\varepsilon) = \frac{B(\varepsilon/2t)}{t}$$

where  $B(x)$  is defined by

$$B(x) = \int_{|\eta| < \pi} \frac{d^3 \vec{\eta}}{(2\pi)^3} \delta \left( x + \sum_{\alpha=x,y,z} \cos \eta_\alpha \right).$$

Here it is noted that the DOS at  $E_F$  becomes

$$D_0(\varepsilon_F = 0) = \frac{B(0)}{t}.$$

**Hint:** To find the expression in the plane-wave-state representation, one may use the following transformation:

$$c_{i\sigma} = \sum_{\mathbf{k} \in \text{B.Z.}} e^{-i\mathbf{k} \cdot \mathbf{R}_i} c_{\mathbf{k}\sigma}$$

and the identity

$$\sum_i e^{i(\mathbf{k}-\mathbf{k}') \cdot \mathbf{R}_i} = \delta_{\mathbf{k},\mathbf{k}'}$$

where  $\mathbf{k}$  and  $\mathbf{k}'$  belong to the first Brillouin zone.

b) **Stoner Theory:**

When  $U$  is small, one can take the Fermi sphere as a *starting* point in considering the effect of Coulomb interaction. Thus, one can rewrite the Eq. (1) by

$$\mathcal{H} = \sum_{\mathbf{k}\sigma} \varepsilon_{0\mathbf{k}} c_{\mathbf{k}\sigma}^\dagger c_{\mathbf{k}\sigma} + U \sum_{\mathbf{q}} n_{-\mathbf{q}\uparrow} n_{\mathbf{q}\downarrow} \quad (3)$$

where  $n_{\mathbf{q}\sigma} = \sum_{\mathbf{k} \in \text{B.Z.}} c_{\mathbf{k}\sigma}^\dagger c_{\mathbf{k}-\mathbf{q},\sigma}$ .

Assuming a *mean field* solution for the ferromagnetic (FM) ground state,

$$\langle n_{\mathbf{q}\sigma} \rangle = \delta_{\mathbf{q},0} \bar{n}_\sigma, \quad (4)$$

show that the *mean field* Hamiltonian can be written by

$$\bar{\mathcal{H}} = \sum_{\mathbf{k}\sigma} \bar{\varepsilon}_{\mathbf{k}\sigma} c_{\mathbf{k}\sigma}^\dagger c_{\mathbf{k}\sigma}$$

with the *spin-dependent* energy term

$$\bar{\varepsilon}_{\mathbf{k}\sigma} = \varepsilon_{0\mathbf{k}} + U \bar{n}_{-\sigma}.$$

Based on the Stoner theory of ferromagnetism, one expects that the ground state of this system changes from paramagnetic to ferromagnetic at a critical value of  $a = a_c$ . Find  $a_c$  in terms of  $A$ ,  $B(0)$ ,  $\gamma$ , and  $U$ , and **DISCUSS** the energetics of the ground states as a function of  $a$ .

c) **Superexchange Interaction:**

Now let us consider the ground state of Eq. (1) in the limit of  $t \rightarrow 0$ , i.e.,  $a \rightarrow \infty$ . An appropriate starting Hamiltonian becomes

$$\mathcal{H}_U = \sum_i n_{i\uparrow} n_{i\downarrow}$$

This limit represents a collection of atoms far apart from each other. Thus the ground state should be a mere superposition of electron in the ground state of each hydrogen atom. Following the *perturbation-theory* arguments for the superexchange mechanism discussed during the class, show that the ground state can be determined by the spin-Hamiltonian

$$\mathcal{H}_{\text{spin}} = J_{\text{AF}} \sum_{\langle ij \rangle} \mathbf{S}_i \cdot \mathbf{S}_j$$

with the antiferromagnetic (AF) exchange interaction, i.e.,  $J_{\text{AF}} > 0$ . Obtain the expression of  $J_{\text{AF}}$  in terms of  $t$  and  $U$ .

d) **Conflicting Results:**

In Question (b), we showed that the ground state becomes FM for  $a > a_c$ , i.e., for small  $t$ , but this result conflicts with the AF ground state as described by the AF spin Hamiltonian in Question (c) even for the same range of  $a$ , i.e.,  $t$ . **DISCUSS** the possible causes of such conflicts, for example, in terms of energetics (total energy) and magnetic instability. (**Hint:** One may repeat the calculations Question (b) for the AF instability.)